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# Forced flame dynamics and self-excited instabilities of laminar premixed flames

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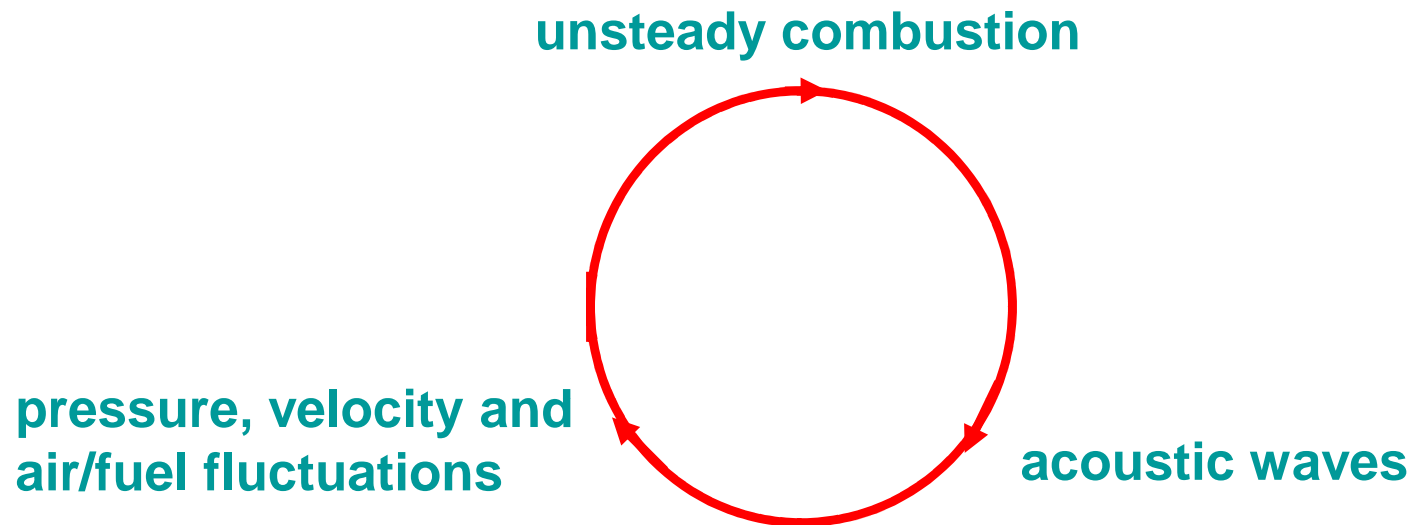
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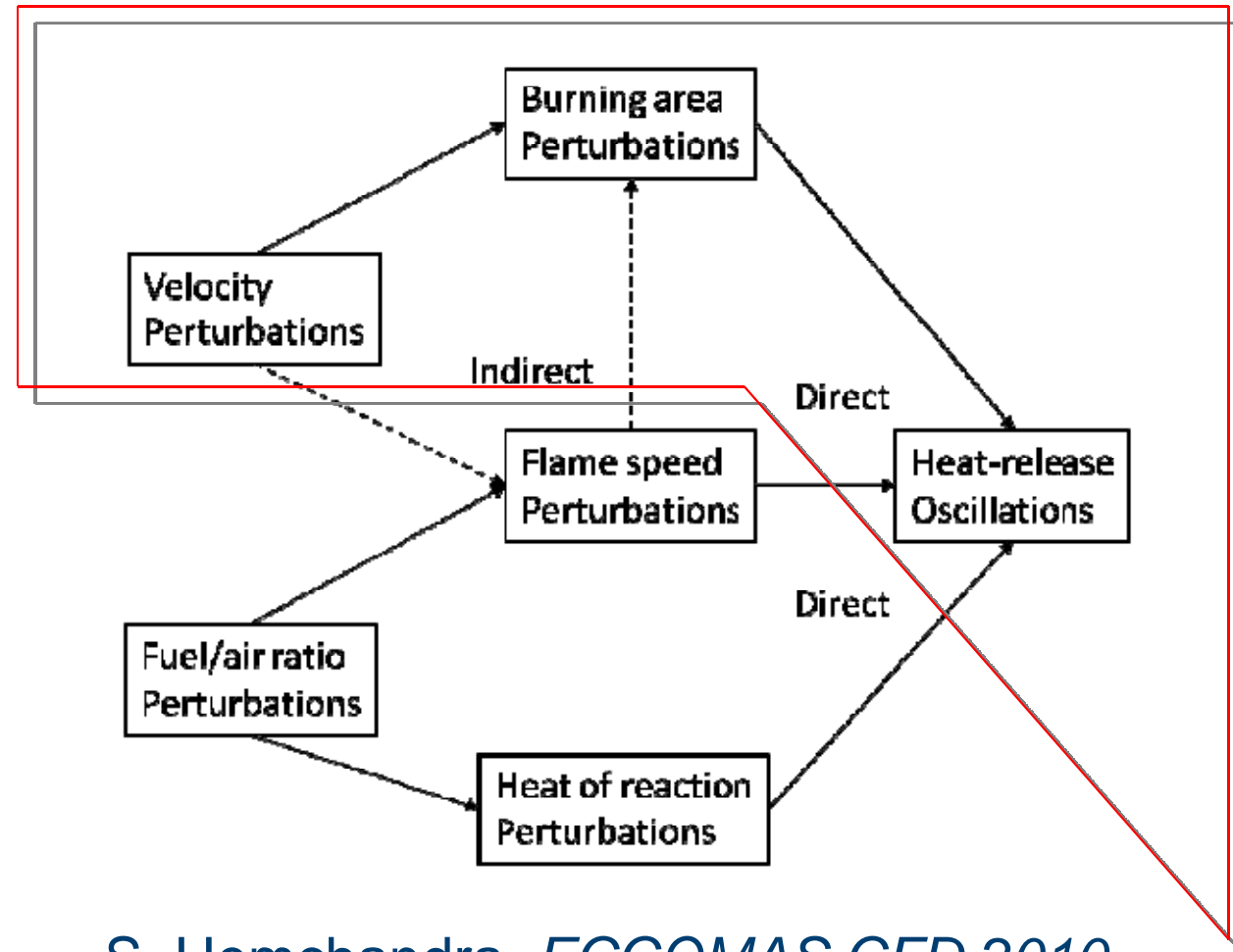
# Thermoacoustic instabilities

Self-excited oscillations due to the coupling shown below



# Unsteady heat release: mechanisms

Premixed flames have unsteady  $q'$  due to  $u'$  only.

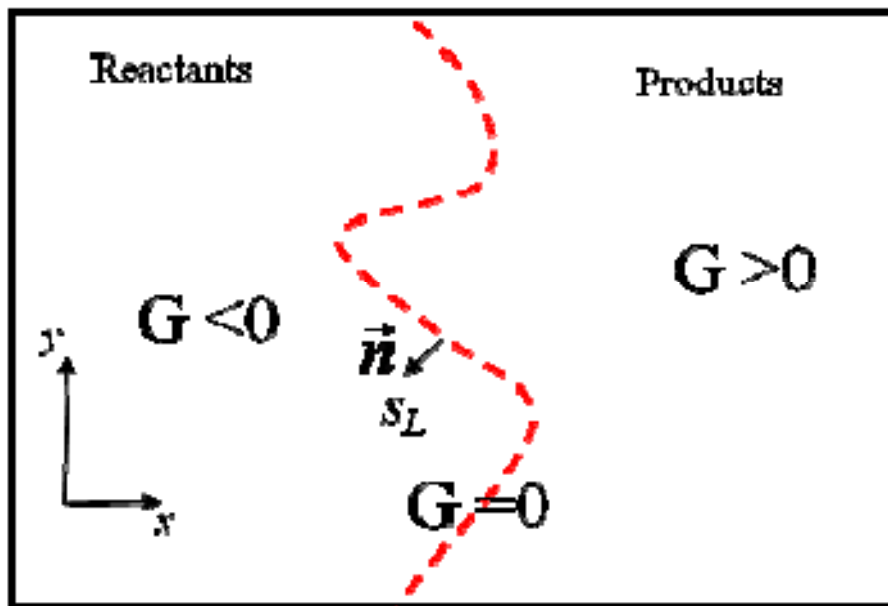


S. Hemchandra, *ECCOMAS CFD 2010*

# Premixed flame model: level set approach

The flame is modelled as a thin moving boundary between reactants and products

Flame represented as a level set of a function  $G(x,y,t)$ .

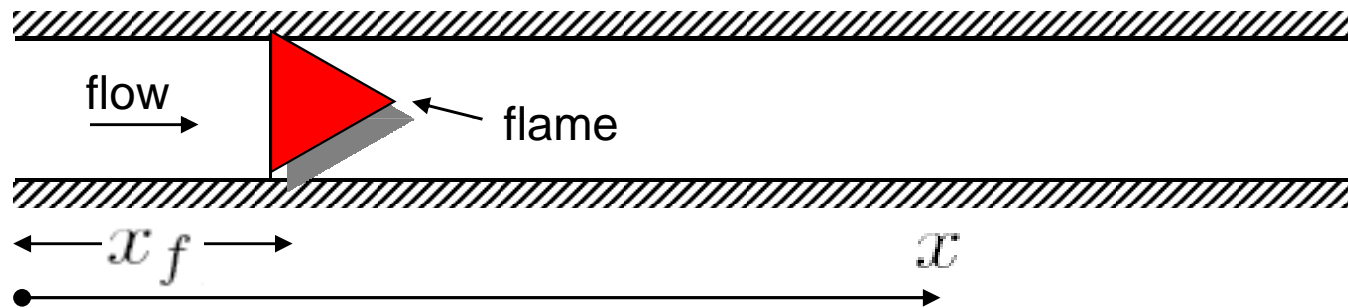


$$\frac{\partial G}{\partial t} + \bar{u}_{flame} \cdot \nabla G = 0$$

$$\frac{\partial G}{\partial t} + \left( \bar{u}(\bar{x}, t) - s_L \frac{\nabla G}{|\nabla G|} \right) \cdot \nabla G = 0$$

# Linear acoustics

Diagram of premixed flame in a tube



Non-dimensional governing equations

$$F_1 \equiv \frac{\partial u}{\partial t} + \frac{\partial p}{\partial x} = 0,$$

$$F_2 \equiv \frac{\partial p}{\partial t} + \frac{\partial u}{\partial x} + \zeta p - \beta \delta_D(x - x_f) = 0$$

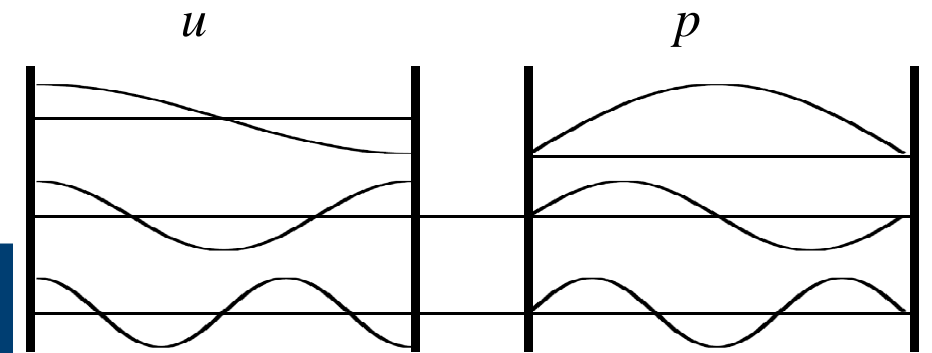
where

$$\beta = \frac{(\gamma - 1) \tilde{Q}_{\tilde{x}_f}}{\mathcal{M} p_0 c_0}$$

Discretization into basis functions

$$u(x, t) = \sum_{j=1}^N \eta_j(t) \cos(j\pi x),$$

$$p(x, t) = -\sum_{j=1}^N \left( \frac{\dot{\eta}_j(t)}{j\pi} \right) \sin(j\pi x),$$



# Flow field: Simplified model of convective disturbances

Acoustic velocity disturbances at the burner lip are convected downstream at a characteristic velocity ( $U_c$ ).

$$\varepsilon = \frac{u'}{u_0}$$

$$u(x, t) = 1 + \varepsilon \cos(kx - \omega t)$$

$$k = \frac{\omega}{u_c} = \frac{\omega}{u_0} \frac{u_0}{u_c} = K \left( \frac{\omega}{u_0} \right)$$

$K$  is a parameter that can be varied in the model

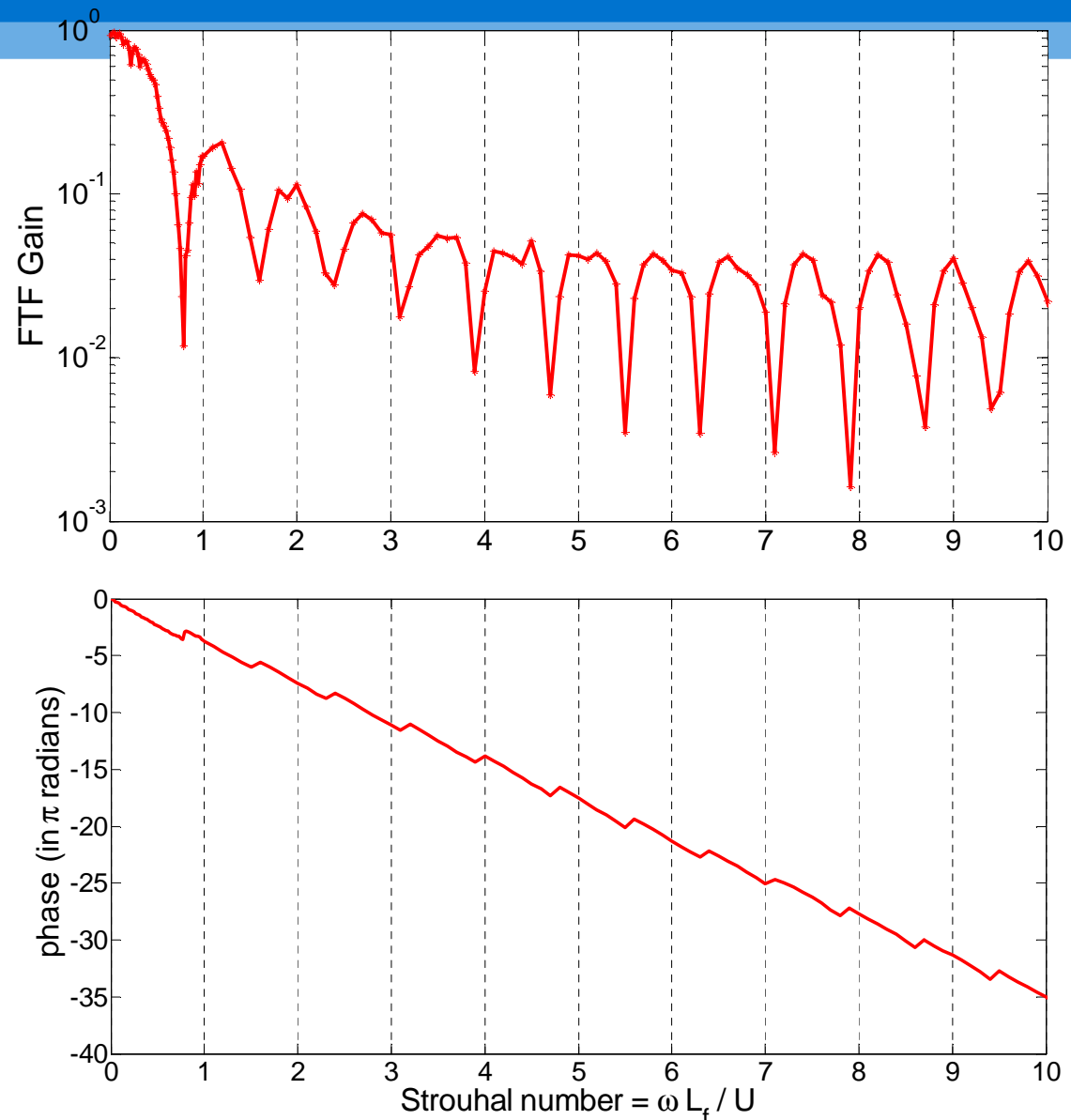


S. Ducruix et. al., *Proc. Comb. Inst.* 2000

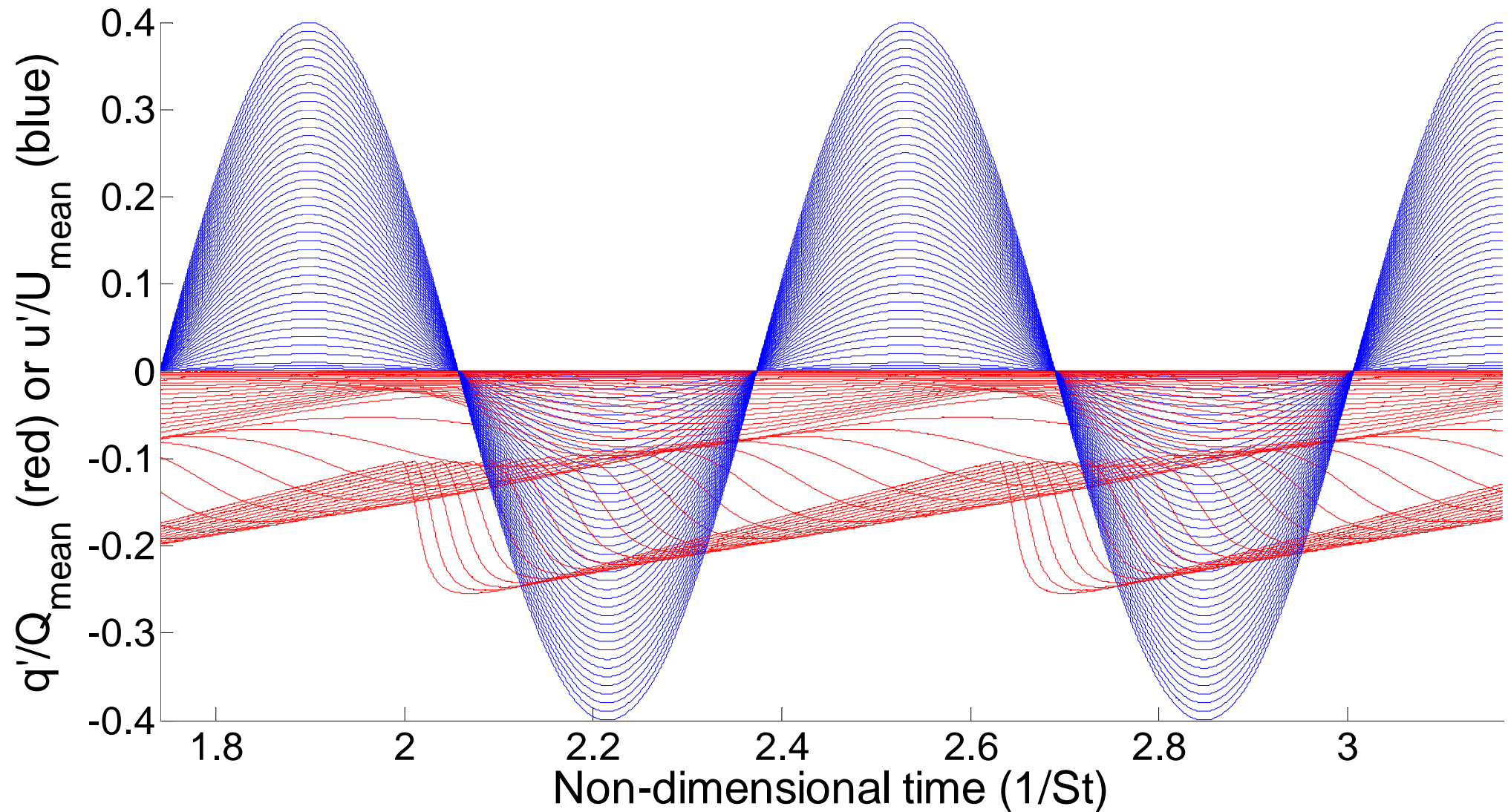
# Linear response: Flame transfer function gain and phase

$$\text{Gain} = \text{Re} \left( \frac{\frac{q'}{Q}}{\frac{u'}{U}} \right)$$

$$\text{Phase} = \text{Angle} \left( \frac{\frac{q'}{Q}}{\frac{u'}{U}} \right)$$



# Nonlinear response of unsteady heat release at $St = 1.6$





# Nonlinear stability criterion for periodic solutions

The stability of periodic solution can be examined as follows (M. P. Juniper)

$$u = r \cos(\omega t + \phi)$$

$$q = \sum_j a_j \cos(j\omega t + \phi_j)$$

$$\oint \frac{dE}{dt} dt = \pi r (a_1 \sin(\omega\tau) - r\zeta\omega)$$

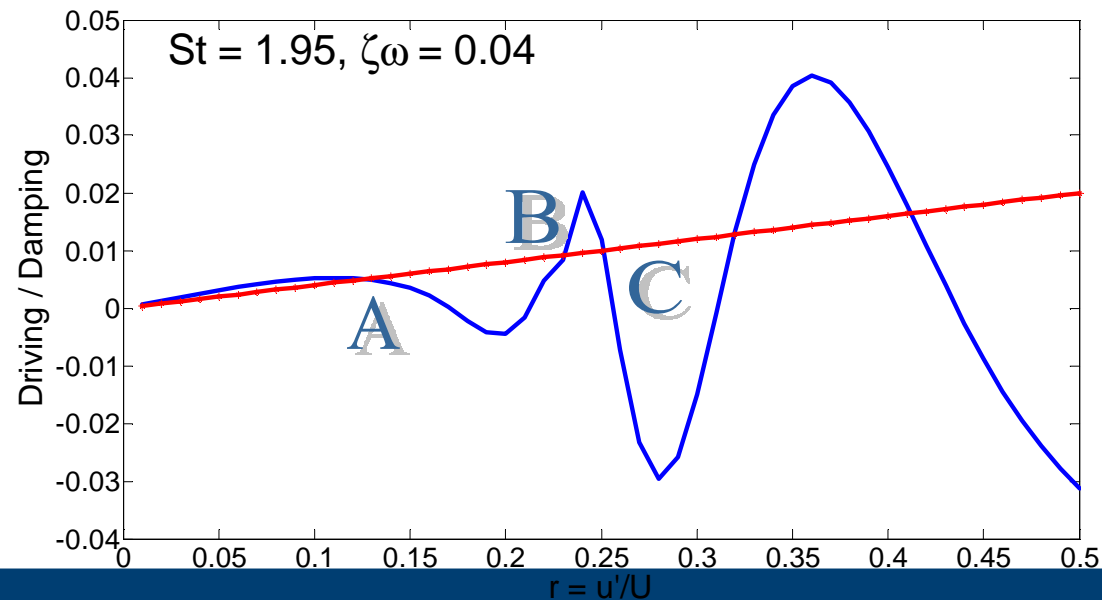
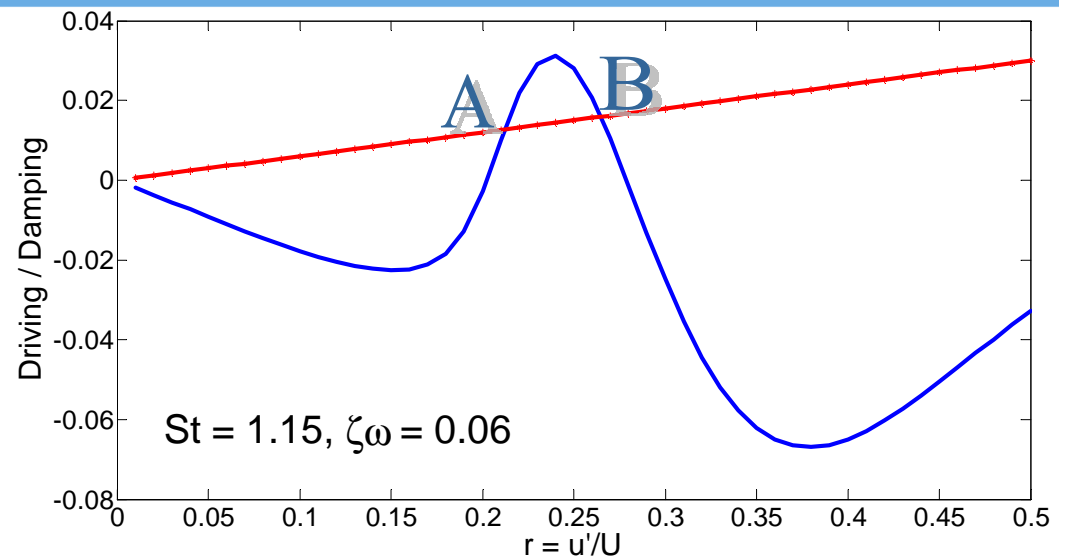
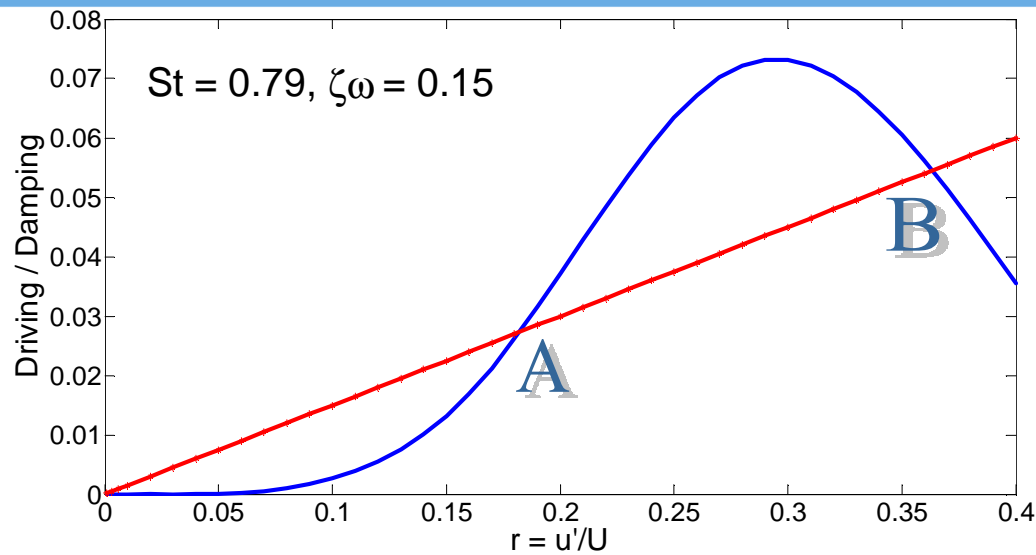
where

$$\tau = \frac{\phi - \phi_1}{\omega}$$

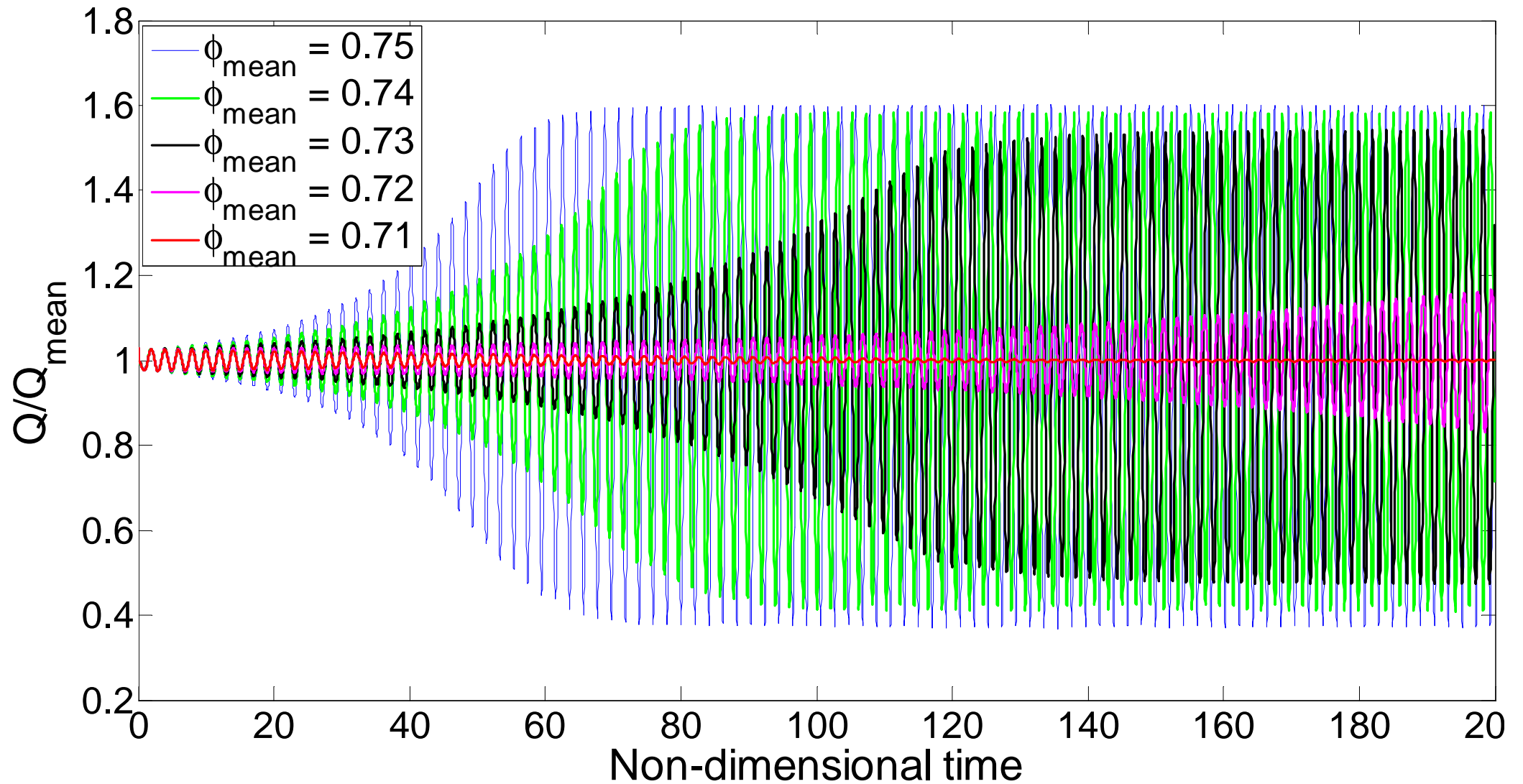
The higher harmonics in the heat release signal do not affect the stability of a periodic solution!

By definition, the integral of  $dE/dt$  goes to zero for a periodic solution. Whether these periodic solutions are stable or unstable can be identified by examining the change in sign of the above integral.

# The premixed flame is susceptible to triggering Both subcritical and supercritical bifurcations are possible



# Self-excited instabilities



# Future work

- Use M.P.Juniper's flow instability tool to predict the wave-speeds of convected disturbances, which is an input to the flame model
- Map out the bifurcation diagram as a function of various parameters (length of tube, flame position, damping factor and flame parameters such as mean equivalence ratio)
- Extend the above analyses to more realistic flame geometries and to stratified flames