

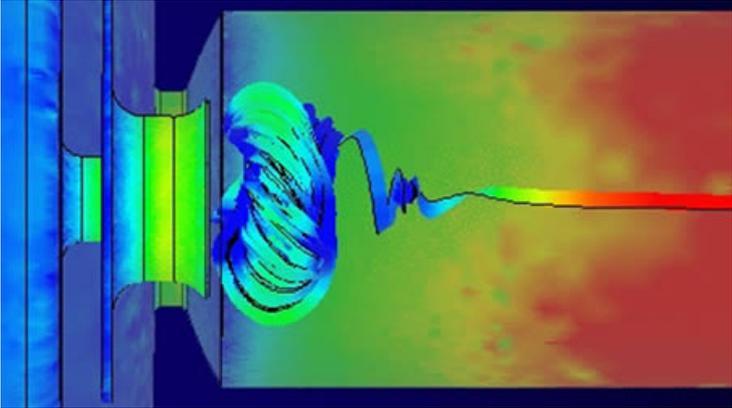
# Sensitivity Analysis of Spiral Vortex Breakdown

Ubaid Ali Qadri, Dhiren Mistry, Matthew Juniper

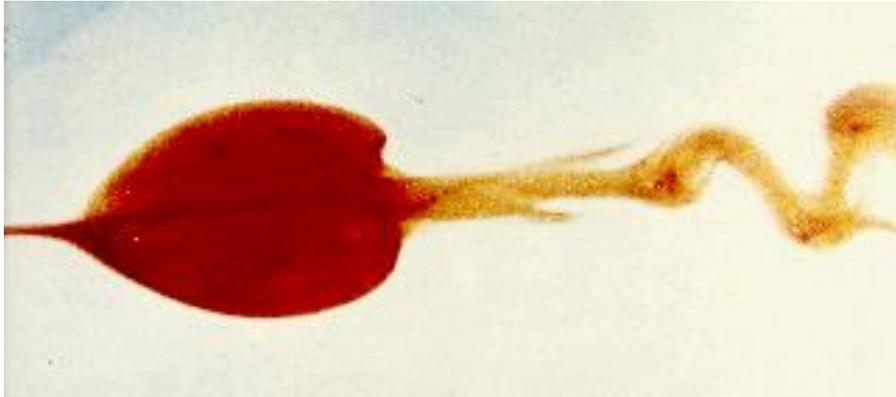
Department of Engineering

We are studying the laminar vortex breakdown bubble because it is a toy model for the recirculating zone in a gas turbine combustion chamber.

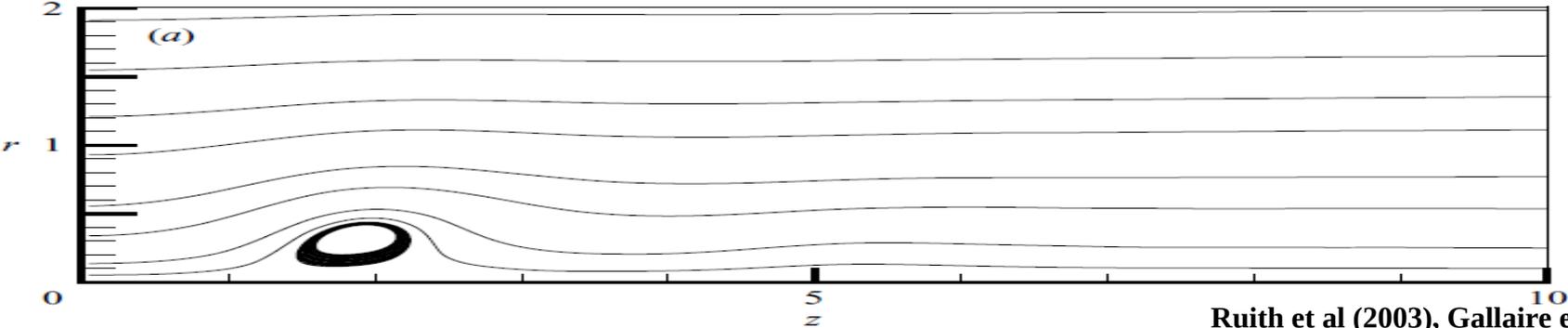
Gas turbine combustion ( $Re \sim 10^6$ )



Vortex breakdown bubble ( $Re \sim 200$ )



Vortex breakdown bubble ( $Sw=1.0, Re = 200$ )

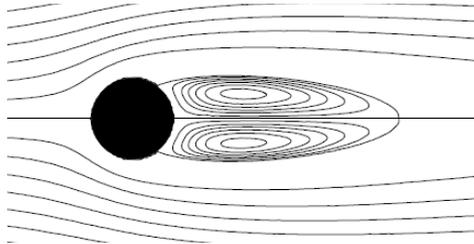


Ruith et al (2003), Gallaire et al (2006)

We use techniques that have been used to study vortex shedding behind a cylinder.

Giannetti & Luchini, JFM (2007) and Hill AIAA (1992)

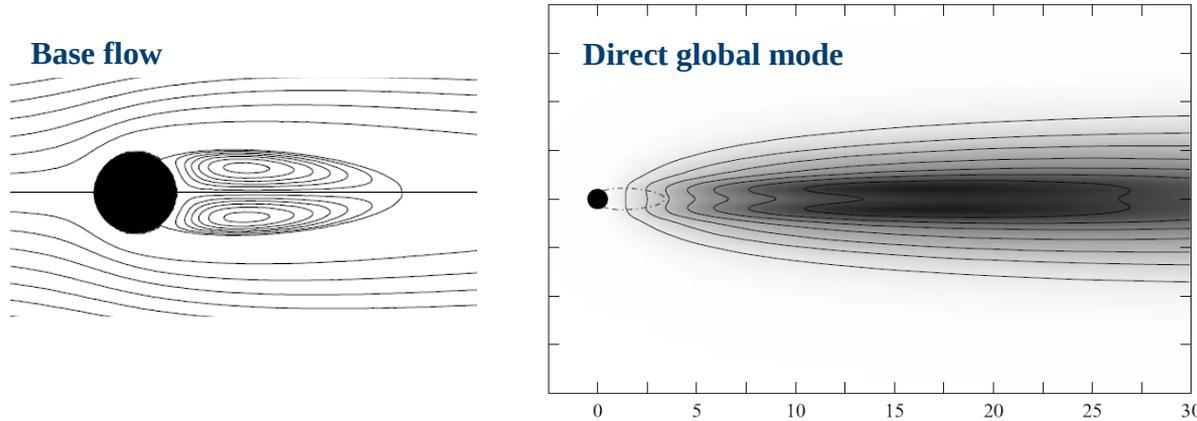
Base flow



Solve Navier-Stokes equations to obtain a steady, axisymmetric solution.

We use techniques that have been used to study vortex shedding behind a cylinder.

Giannetti & Luchini, JFM (2007) and Hill AIAA (1992)



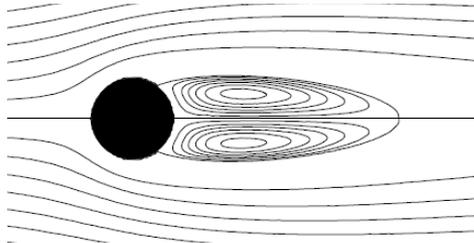
Superpose small perturbations onto steady base flow.  
Solve 2D eigenvalue problem.

$$\frac{dq}{dt} = \mathbf{L}q,$$
$$q(x, r, \theta, t) = \hat{q}(x, r)e^{im\theta + \lambda t}, \quad \lambda \equiv \sigma + i\omega$$
$$\lambda \hat{q} = \mathbf{L}_m \hat{q}.$$

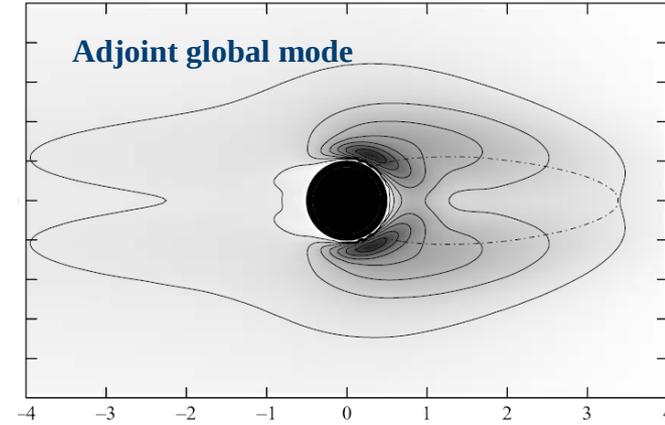
# We use techniques that have been used to study vortex shedding behind a cylinder.

Giannetti & Luchini, JFM (2007) and Hill AIAA (1992)

Base flow



Adjoint global mode



Derive adjoint (conjugate transpose) equations.  
Solve another 2D eigenvalue problem.

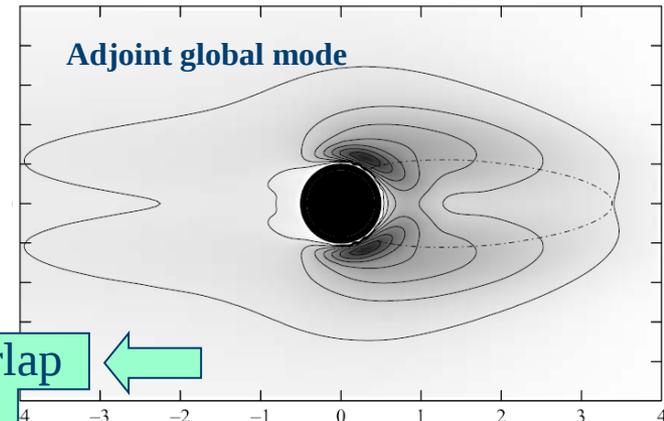
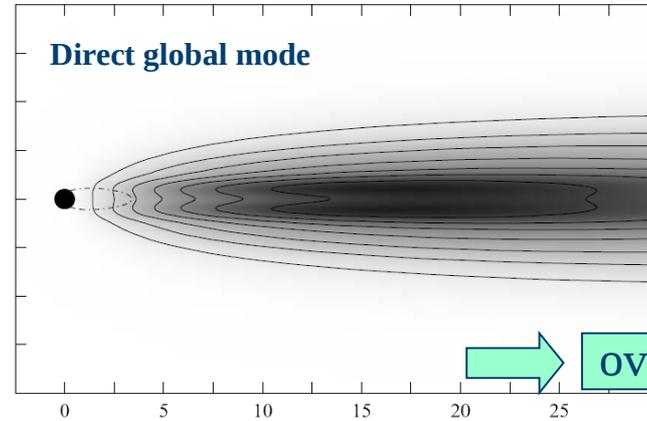
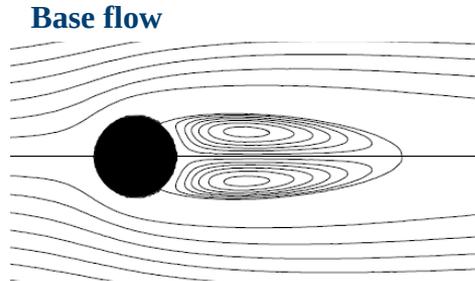
$$-\frac{dq^+}{dt} = \mathbf{L}^+ q^+$$

$$q^+(x, r, \theta, t) = \hat{q}^+(x, r) e^{im\theta - \lambda^* t}$$

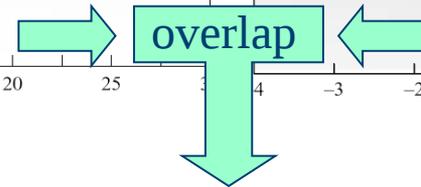
$$\lambda^* \hat{q}^+ = \tilde{\mathbf{L}}_m^+ \hat{q}^+$$

# We use techniques that have been used to study vortex shedding behind a cylinder.

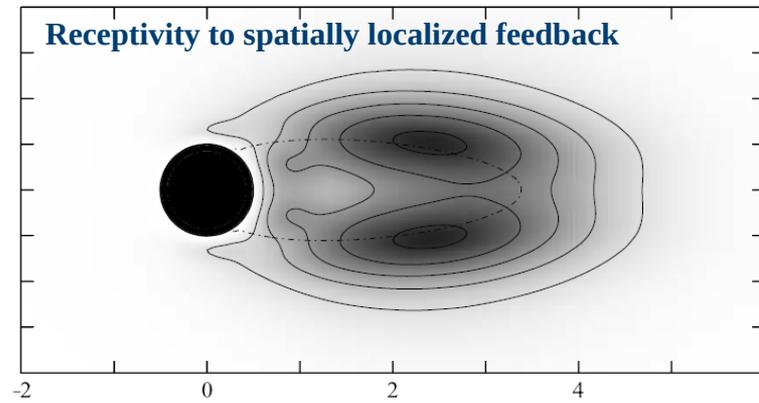
Giannetti & Luchini, JFM (2007) and Hill AIAA (1992)



overlap

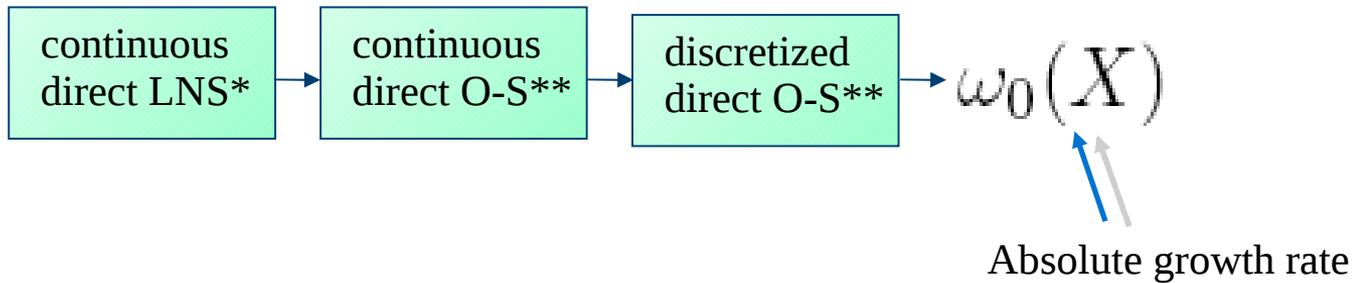
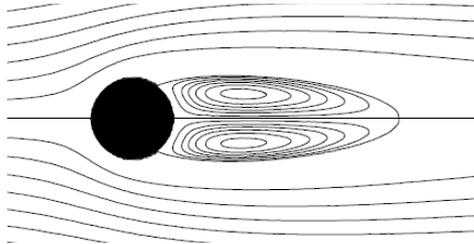
A green box with the word "overlap" inside. Two green arrows point towards the box from the left and right, and a larger green arrow points downwards from the box.

$$|\nabla_G \lambda|_{max} = \frac{|\hat{m}^H| |\hat{m}^+|}{\langle \hat{m}^+, \hat{m} \rangle}.$$



The absolute growth rate is a useful diagnostic result. It is obtained from a local stability analysis using INSTAFLOW.

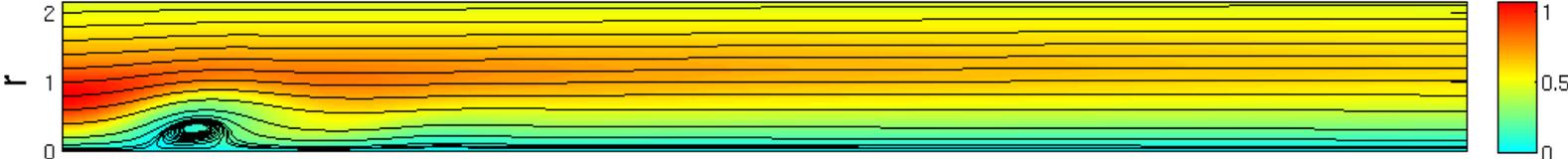
Base flow



\* LNS = Linearized Navier-Stokes equations

\*\* O-S = Orr-Sommerfeld equation

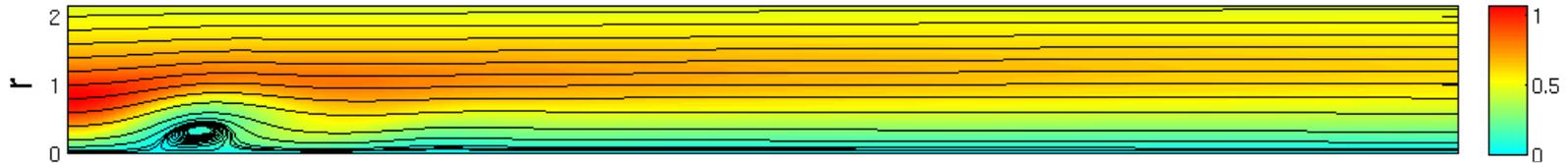
We checked the results for the linear direct global mode against previously-published results for the nonlinear global mode.



**Direct eigenvalues**

$S_W=1.0$	Xmax	Rmax	Sx	Sr	Growth rate	Frequency
Fully 3D DNS (Ruith)	20.0	10.0	193	61	0.0359	1.1800
Local nonlinear analysis (Gallaire)	-	-	-	-	-	1.2200
M1	20.0	8.0	513	127	0.035214	1.165476

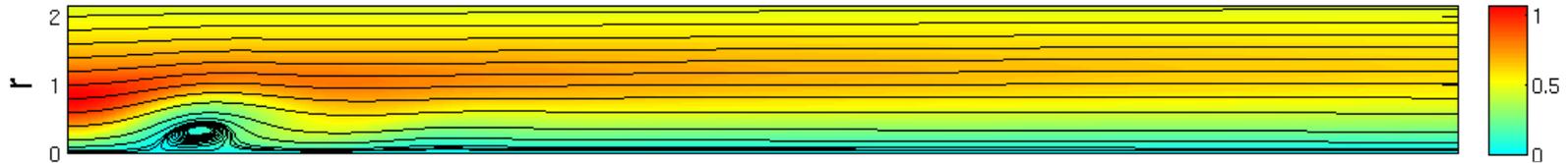
We checked the effect of the grid resolution and size of the domain.



### Direct eigenvalues

$Sw=1.0$	Xmax	Rmax	Sx	Sr	Growth rate	Frequency
Fully 3D DNS (Ruith)	20.0	10.0	193	61	0.0359	1.1800
Local nonlinear analysis (Gallaire)	-	-	-	-	-	1.2200
M1	20.0	8.0	513	127	0.035214	1.165476
M2	20.0	8.0	257	127	0.035177	1.165453
M3	25.4	10.4	257	127	0.034348	1.162470

We checked the results for the adjoint global mode against the direct global mode.



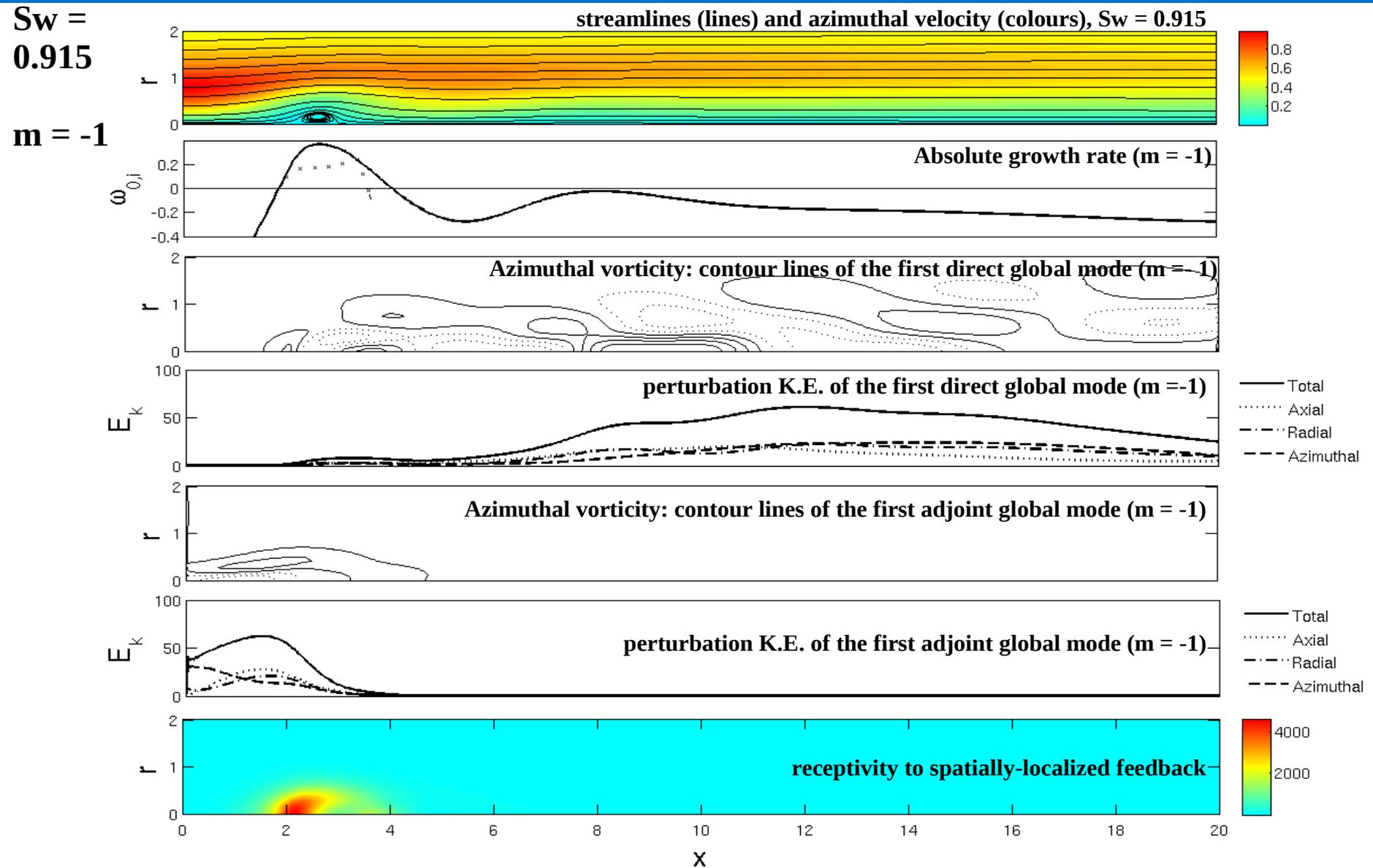
### Direct eigenvalues

$Sw=1.0$	Xmax	Rmax	Sx	Sr	Growth rate	Frequency
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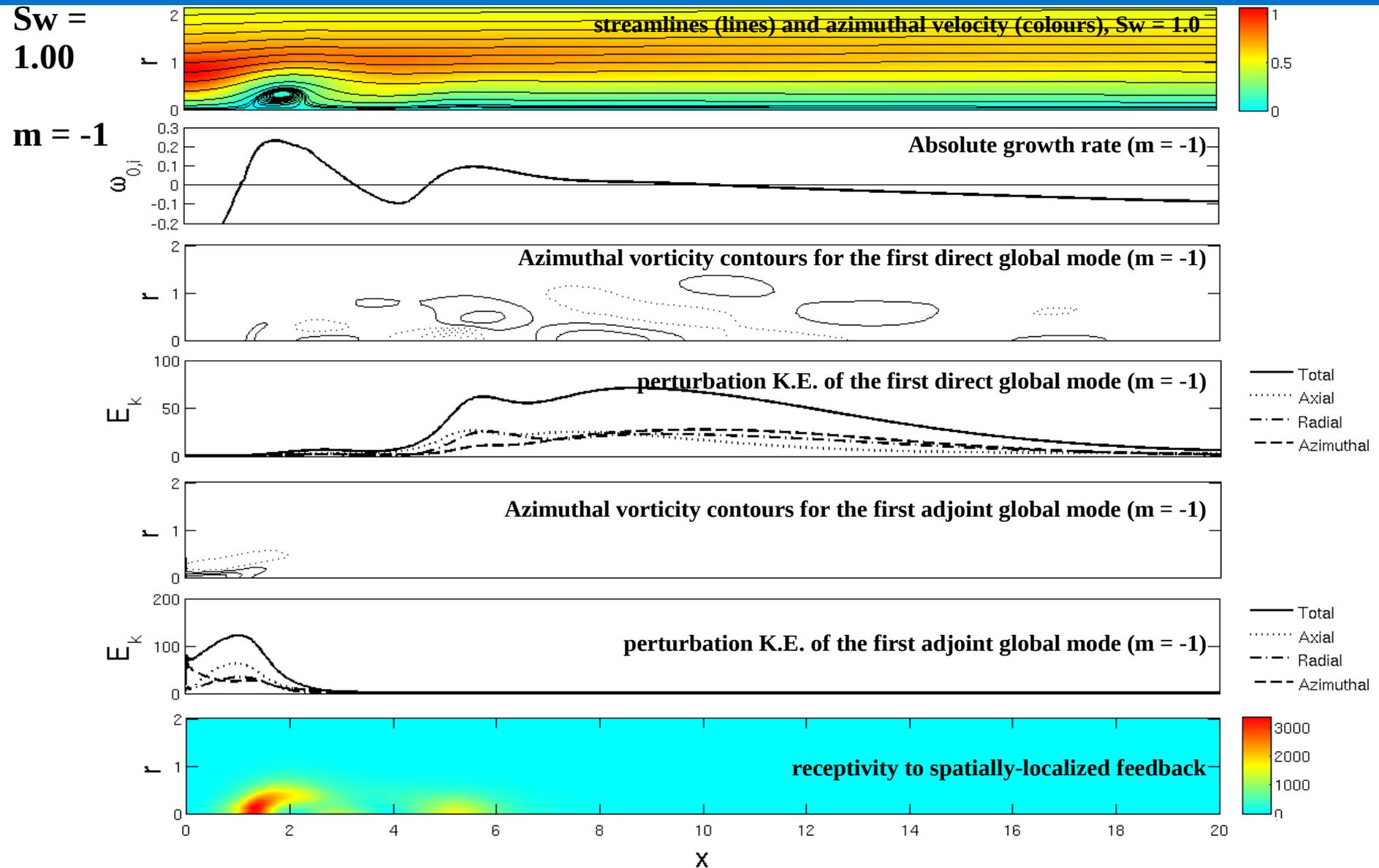
### Adjoint eigenvalues

	Growth rate	Frequency	Absolute discrepancy	Relative discrepancy (%)
M1	0.037663	-1.165434	0.002449	0.210064

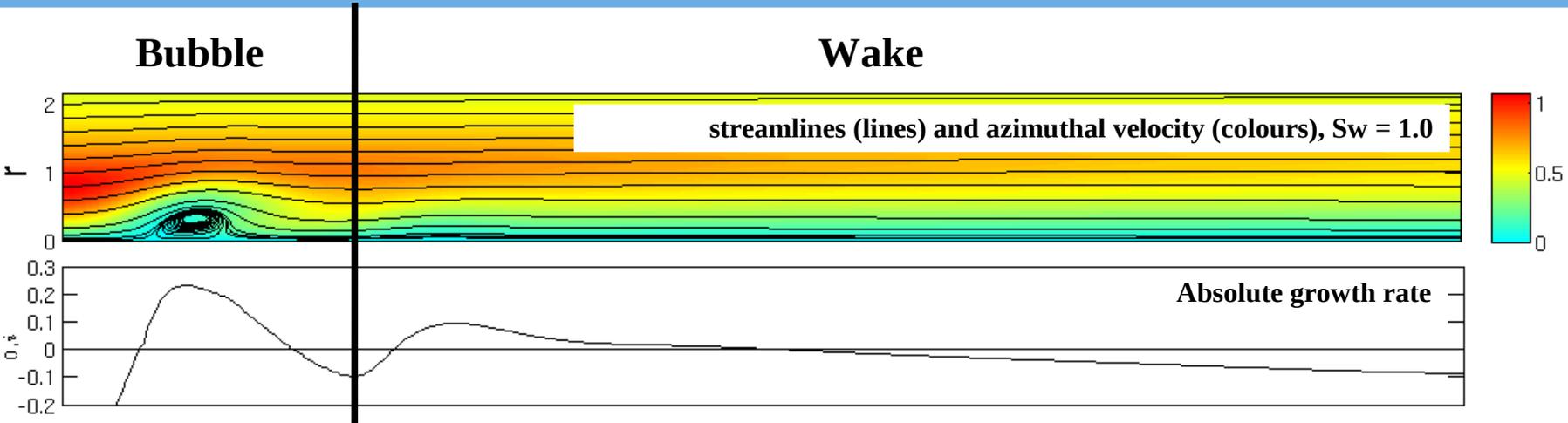
At  $Sw=0.915$ , the flow is just globally unstable for  $m=-1$ . The recirculation bubble is absolutely unstable. The overlap of direct and adjoint modes shows that the wavemaker region is in the recirculation bubble.



At  $Sw=1.0$ , there are two regions of absolute instability. The overlap of direct and adjoint modes shows that the wavemaker region is still predominantly in the recirculation bubble.

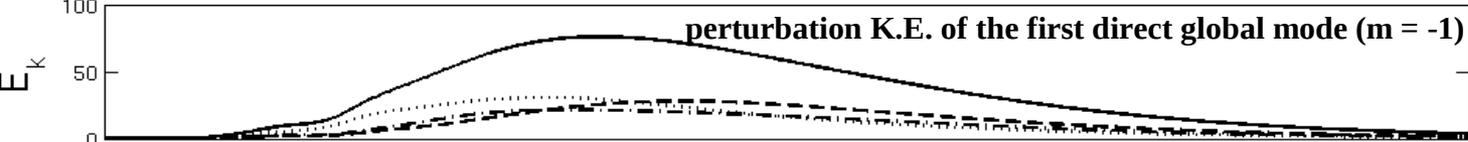
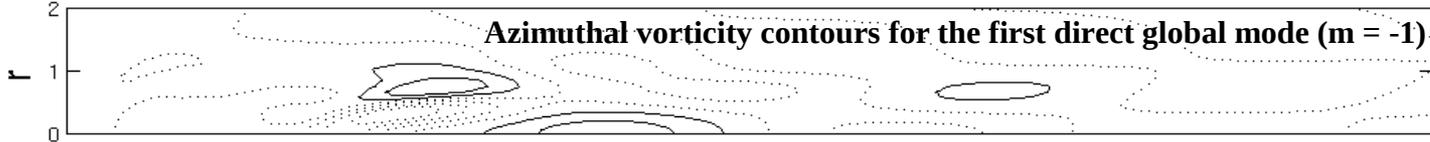
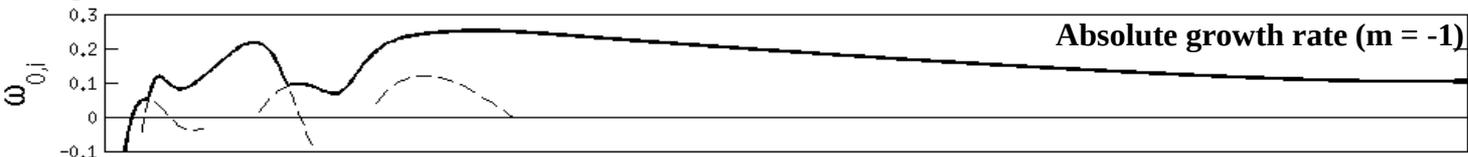
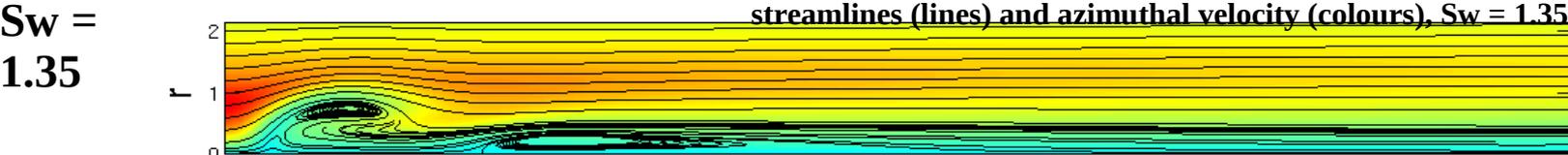


We calculated global modes of the bubble and wake separately. The frequency of the global mode seems to be driven by the bubble but the growth rate seems to be enhanced by the nearby marginally-stable wake. It is like a coupled oscillator.

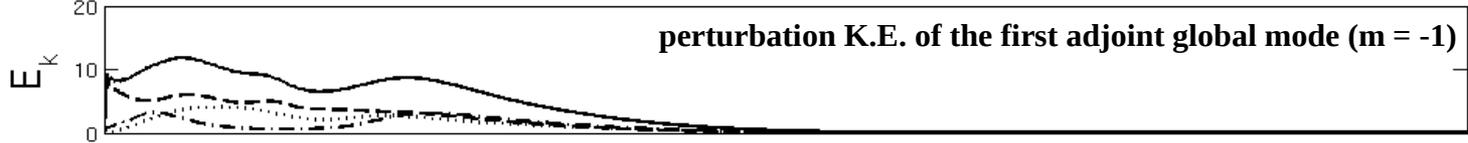
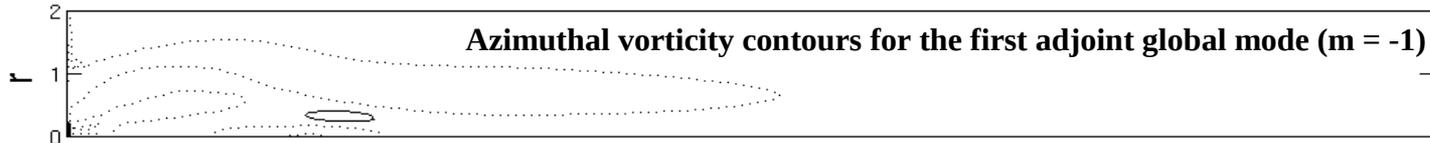


	Growth rate	Frequency
<b>Complete global</b>	<b>0.0352</b>	<b>1.1655</b>
<b>Bubble</b>	<b>0.0132</b>	<b>1.1698</b>
<b>Wake</b>	<b>-0.0475</b>	<b>1.0852</b>

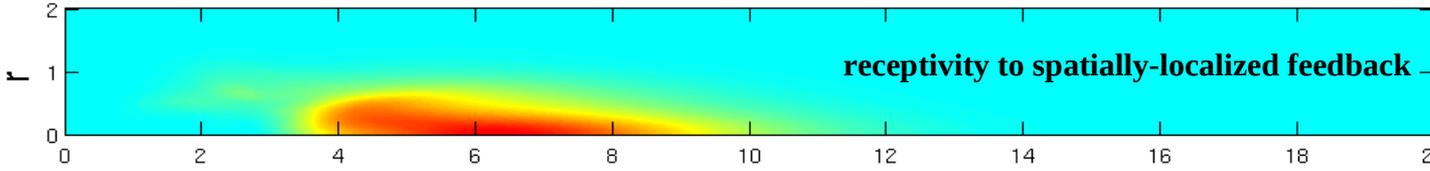
As the swirl increases, a second recirculation bubble forms in the wake. The wavemaker region moves into the wake bubble.



- Total
- ..... Axial
- - - Radial
- - - Azimuthal

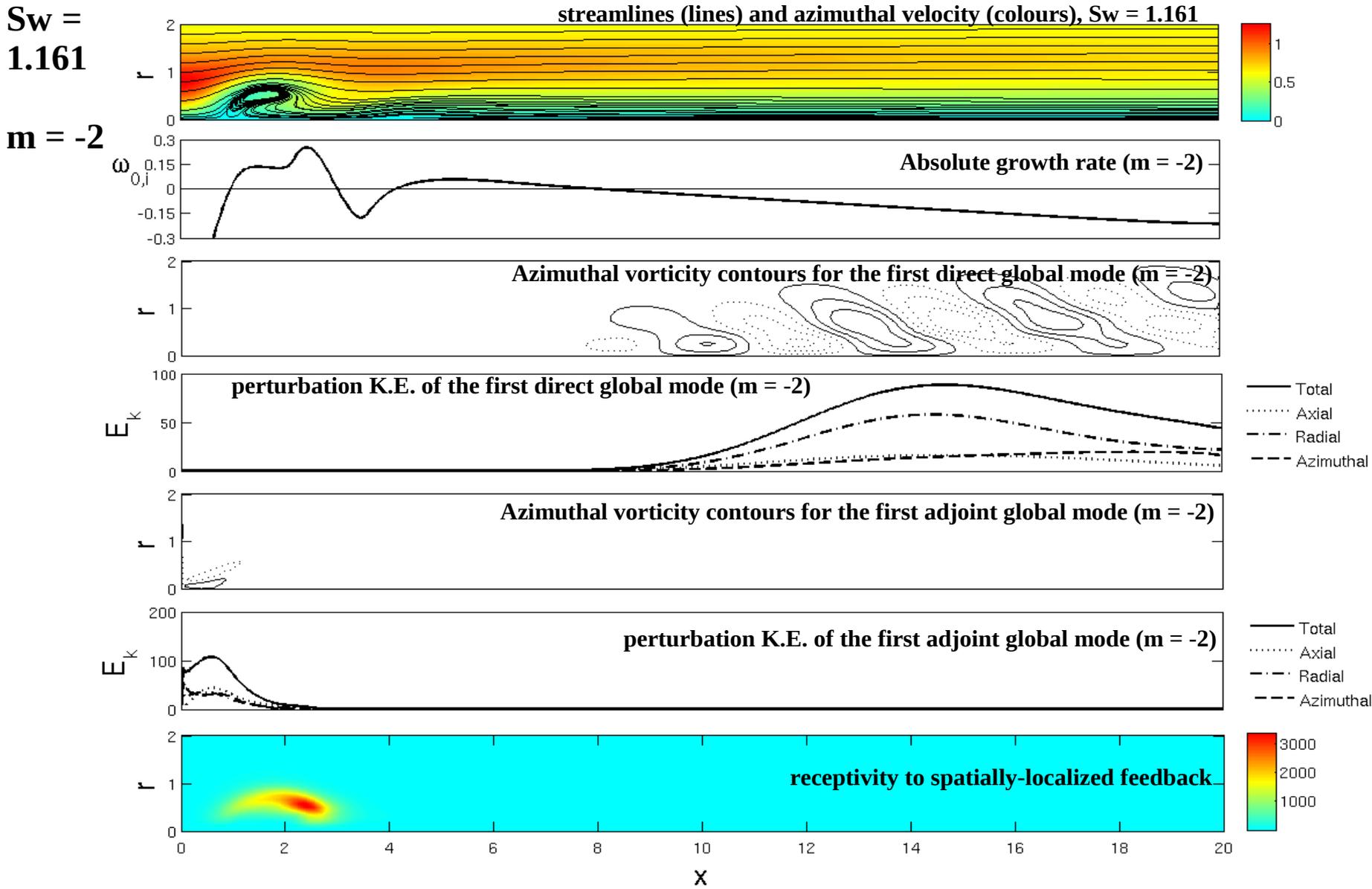


- Total
- ..... Axial
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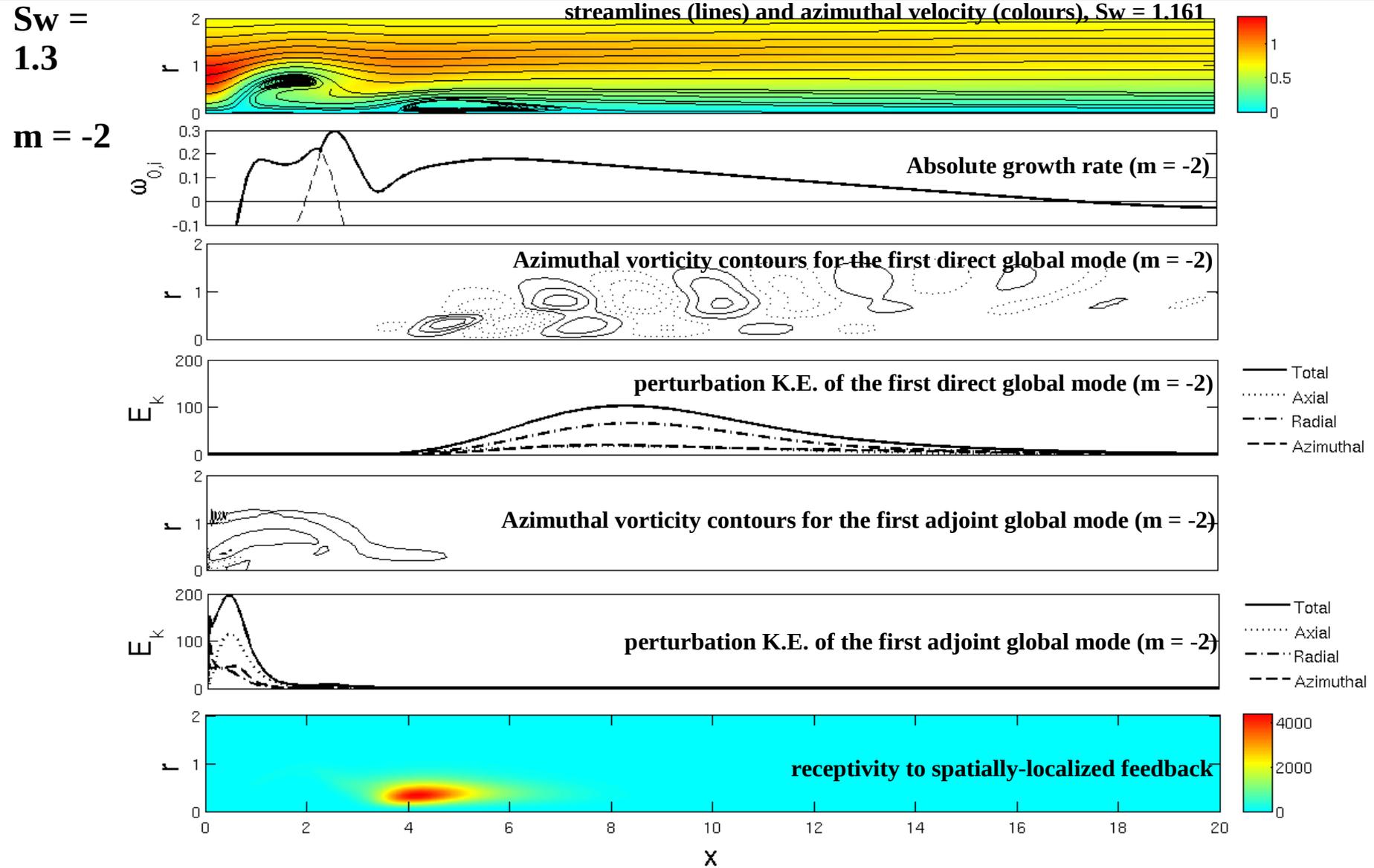


X

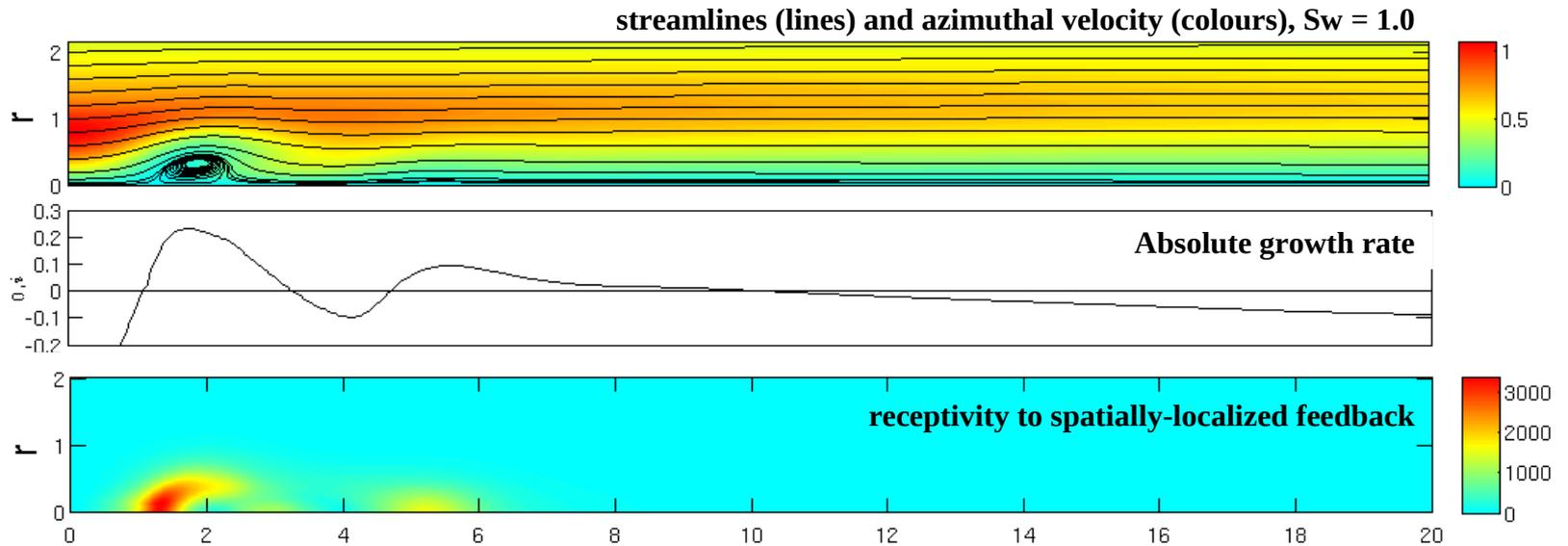
The  $m = -2$  mode becomes unstable at  $Sw = 1.161$ . The wavemaker region is in the recirculation bubble.



As the swirl increases, the wavemaker region moves into the recirculating wake, as for the  $m = -1$  mode.

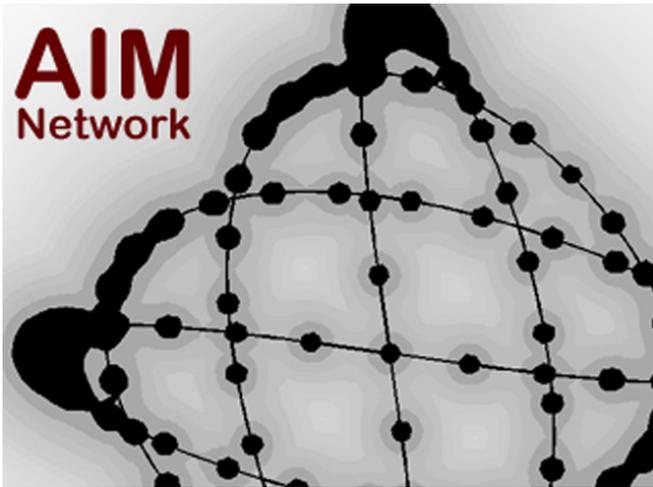


**1. At low swirls, the wavemaker region is in the bubble, not the wake.**



**2. As the swirl increases, the wavemaker region moves into the wake.**

**3. These effects are seen for both the  $m = -1$  and  $m = -2$  modes.**



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*Monday 22nd August to  
Saturday 27th August 2011,*

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